TFY4210, Quantum theory of many-particle systems, 2016: Tutorial 9

In this tutorial set, all problems concern fermions.

1. Relationship between $G^{<}(\nu,\omega)$ and the spectral function $A(\nu,\omega)$.

(a) Starting from the definition of the "lesser" Green's function, show that it can be written

$$G^{<}(\nu;t,t') = \frac{i}{Z} \sum_{n,m} e^{-\beta E_m} e^{i(E_n - E_m)(t - t')} |\langle m|c_{\nu}^{\dagger}|n\rangle|^2. \tag{1}$$

and that its Fourier transform is

$$G^{<}(\nu,\omega) = \frac{2\pi i}{Z} \sum_{n,m} e^{-\beta E_m} |\langle m|c_{\nu}^{\dagger}|n\rangle|^2 \delta(\omega + E_n - E_m). \tag{2}$$

(b) Starting from Eq. (53) in the lecture notes, show that $A(\nu,\omega)$ can be rewritten as

$$A(\nu,\omega) = (1 + e^{\beta\omega}) \frac{1}{Z} \sum_{n,m} |\langle m|c_{\nu}^{\dagger}|n\rangle|^2 e^{-\beta E_m} \delta(\omega + E_n - E_m). \tag{3}$$

(c) Show that

$$-iG^{<}(\nu,\omega) = 2\pi A(\nu,\omega)n_F(\omega) \tag{4}$$

which is Eq. (56) in the lecture notes. (Eq. (55) can be proved in a similar way.)

2. An alternative form of the Lehmann representation.

Show that [Eq. (54) in the lecture notes]

$$G^{R}(\nu,\omega) = \int_{-\infty}^{\infty} d\omega' \, \frac{A(\nu,\omega')}{\omega - \omega' + i\eta}.$$
 (5)

3. Calculating $G^R(\nu,t)$ from $G^R(\nu,\omega)$ by contour integration.

In the lectures we calculated $G^R(\nu, \omega)$ from a knowledge of $G^R(\nu, t)$. But suppose that you instead know $G^R(\nu, \omega)$ and want to calculate $G^R(\nu, t)$ from it. Starting from the Lehmann representation for $G^R(\nu, \omega)$ [Eq. (50) in the lecture notes], calculate

$$G^{R}(\nu,t) = \frac{1}{2\pi} \int_{-\infty}^{\infty} d\omega e^{-i\omega t} G^{R}(\nu,\omega)$$
 (6)

by using **contour integration**: Consider the cases t > 0 and t < 0 separately. For each case, identify which half-plane the contour must be closed in. By using the residue theorem to evaluate the resulting contour integrals, show that

$$G^{R}(\nu,t) = -i\theta(t)\frac{1}{Z}\sum_{n,m} \left(e^{-\beta E_n} + e^{-\beta E_m}\right)e^{i(E_n - E_m)t}|\langle m|c_{\nu}^{\dagger}|n\rangle|^2 \tag{7}$$

which is Eq. (48) in the lecture notes.

4. Fermi liquids.

Consider the following simplified model for the spectral function of a Fermi liquid for wavevectors k in the vicinity of the Fermi surface:

$$A(\mathbf{k}\sigma,\omega) = Z\delta(\omega - \xi_{\mathbf{k}}) + (1 - Z) \frac{\theta(W - |\omega|)}{2W}.$$
 (8)

Here Z and W are constants (real and positive, with $Z \leq 1$), and $\theta(x)$ is the Heaviside (step) function.

- (a) Briefly compare and contrast this expression with the spectral function of a Fermi liquid as discussed in the lecture notes [Eq. (60)].
- (b) Show that (8) satisfies the sum rule

$$\int_{-\infty}^{\infty} d\omega \, A(\mathbf{k}\sigma, \omega) = 1. \tag{9}$$

(c) Show that at zero temperature the momentum distribution function $\langle c_{\mathbf{k}\sigma}^{\dagger} c_{\mathbf{k}\sigma} \rangle$ of the Fermi liquid described by (8) has a jump of magnitude Z as $k = |\mathbf{k}|$ crosses the Fermi surface.