TFY4210, Quantum theory of many-particle systems, 2014: Tutorial 2

1. Explicit connection between first and second quantization.

An explicit connection between the basis states $|n\rangle \equiv |n_1, n_2, \ldots\rangle$ with $\sum_{\nu} n_{\nu} = N$ in second quantization and the basis wavefunctions $\Phi_n(x_1, \ldots, x_N)$ in first quantization can be established. Let

$$|x_1, x_2, \dots, x_N\rangle \equiv \frac{1}{\sqrt{N!}} \hat{\psi}^{\dagger}(x_1) \dots \hat{\psi}^{\dagger}(x_N) |0\rangle.$$
 (1)

Then the (correctly normalized) basis wavefunctions are given by

$$\Phi_n(x_1, \dots, x_N) = \langle x_1, \dots, x_N | n \rangle. \tag{2}$$

As an example, consider a fermionic 2-particle basis state $|\ldots, 1_{\mu}, \ldots, 1_{\nu}, \ldots\rangle$ in which the single-particle states μ and ν are occupied and all others are empty. Evaluate the rhs of (2) to show that the wavefunction is indeed the correct Slater determinant,

$$\Phi(x_1, x_2) = \frac{1}{\sqrt{2!}} \begin{vmatrix} \phi_{\mu}(x_1) & \phi_{\mu}(x_2) \\ \phi_{\nu}(x_1) & \phi_{\nu}(x_2) \end{vmatrix}.$$
 (3)

2. Density operators.

The density operator $\hat{\rho}(x)$ is in first quantization given by $\hat{\rho}(x) = \sum_{i=1}^{N} \delta(x - x_i)$.

(a) Show that in second quantization,

$$\hat{\rho}(x) = \hat{\psi}^{\dagger}(x)\hat{\psi}(x). \tag{4}$$

(b) Using second quantization, show that

$$\int dx \ \hat{\rho}(x) = \hat{N} \tag{5}$$

where \hat{N} is the total number operator.

3. Proof of the second-quantized representation of two-particle operators.

In first quantization a two-particle operator \hat{H}_I takes the form

$$\hat{H}_I = \frac{1}{2} \sum_{\substack{i,j=1\\i \neq j}}^{N} v(x_i, x_j).$$
 (6)

(a) Show that this can be rewritten (with $\hat{\rho}(x)$ given by its first-quantization expression)

$$\hat{H}_I = \frac{1}{2} \left[\int dx \int dx' \ v(x, x') \hat{\rho}(x) \hat{\rho}(x') - \int dx \ v(x, x) \hat{\rho}(x) \right]. \tag{7}$$

(b) Show that in second quantization \hat{H}_I can be written (for both fermionic and bosonic systems)

$$\hat{H}_I = \frac{1}{2} \int dx \int dx' \ v(x, x') \ \hat{\psi}^{\dagger}(x) \hat{\psi}^{\dagger}(x') \hat{\psi}(x') \hat{\psi}(x). \tag{8}$$

(c) Use this to show that the second-quantized representation of H_I , expressed using the arbitrary basis $\{|\alpha\rangle\}$ for single-particle states, is given by

$$\hat{H}_{I} = \frac{1}{2} \sum_{\alpha,\beta,\gamma,\delta} \left(\int dx \int dx' \, \phi_{\alpha}^{*}(x) \phi_{\beta}^{*}(x') v(x,x') \phi_{\delta}(x') \phi_{\gamma}(x) \right) \hat{c}_{\alpha}^{\dagger} \hat{c}_{\beta}^{\dagger} \hat{c}_{\delta} \hat{c}_{\gamma}, \tag{9}$$

which is exactly of the form claimed in the lectures (the expression enclosed in parentheses is the matrix element denoted $\langle \alpha \beta | v | \gamma \delta \rangle$ in the lectures).