## TFY4210, Quantum theory of many-particle systems, 2015: Tutorial 7

## 1. Bogoliubov transformation for bosons.

Consider the Hamiltonian

$$H = \varepsilon(a_1^{\dagger} a_1 + a_2^{\dagger} a_2) + \Delta(a_1 a_2 + \text{h.c.})$$
 (1)

where  $a_1$  and  $a_2$  are bosonic operators satisfying canonical commutation relations  $[a_i, a_j^{\dagger}] = \delta_{ij}$  etc. (i, j = 1, 2). Assume that  $\varepsilon$  and  $\Delta$  are positive numbers with  $\varepsilon > \Delta$ .

In order to write the Hamiltonian in diagonal form we do a Bogoliubov transformation to a new set  $b_1$ ,  $b_2$  of bosonic operators. The transformation reads (here u and v are real numbers)

$$a_1 = ub_1 - vb_2^{\dagger}, \tag{2}$$

$$a_2 = ub_2 - vb_1^{\dagger}. \tag{3}$$

(a) Use the requirement that the *b*-operators should also satisfy canonical commutation relations to show (e.g. by just calculating one selected commutator) that

$$u^2 - v^2 = 1. (4)$$

This result can be used to write  $u = \cosh \eta$ ,  $v = \sinh \eta$ .

(b) Show that H becomes diagonal in terms of the b-operators provided that  $\eta$  is chosen to satisfy

$$\tanh 2\eta = \frac{\Delta}{\varepsilon}.\tag{5}$$

Show that with this choice, H can be written

$$H = F(b_1^{\dagger}b_1 + b_2^{\dagger}b_2) + G \tag{6}$$

and give expressions for F and G in terms of  $\varepsilon$  and  $\Delta$ .

- (c) Argue from this result that the ground state  $|\Psi_0\rangle$  can be defined in terms of an equation expressing what happens when an annihilation operator  $b_i$  acts on  $|\Psi_0\rangle$ . What is the ground state energy? Explain your reasoning.
- (d) What is the energy of the lowest excited state(s)? Explain your reasoning.
- (e) Express the annihilation operator  $b_1$  in terms of a-operators.
- (f) The ground state  $|\Psi_0\rangle$  is given by

$$|\Psi_0\rangle = C \exp(-\tanh \eta \ a_1^{\dagger} a_2^{\dagger})|0\rangle$$
 (7)

where  $|0\rangle$  is the vacuum state of the a-operators, i.e.  $a_1|0\rangle = a_2|0\rangle = 0$ . The constant C is just a normalization factor and can be neglected in the following.

Verify Eq. (7) for the ground state by showing that it satisfies the defining equation for  $|\Psi_0\rangle$  discussed in (c) (it's sufficient that you only consider the equation that involves  $b_1$ ). [Hint: The Baker-Hausdorff theorem

$$e^{-Q}Pe^{Q} = P + [P,Q] + \frac{1}{2!}[[P,Q],Q] + \frac{1}{3!}[[[P,Q],Q],Q] + \dots$$
 (8)

may be useful.]

## 2. Ferromagnetic Heisenberg model with a spin anisotropy revisited.

Consider again the ferromagnetic Heisenberg model with a spin anisotropy studied in Problem 2 in Tutorial 6, where for  $D \geq 0$  it was found that the energy gap  $\Delta = 2SD$ . Thus the system is gapped for D > 0 and gapless for D = 0. Discuss this property in light of Goldstone's theorem and the symmetries of the model and its ground states.

## 3. The total $S^z$ operator for the antiferromagnetic Heisenberg model.

Consider the operator  $S^z = \sum_j S_j^z$  representing the total spin in the z direction (the sum is over all N sites in the system).

(a) Use spin-wave theory to show that for the antiferromagnetic Heisenberg model, one can express  $S^z$  as

$$S^{z} = \sum_{k} (\beta_{k}^{\dagger} \beta_{k} - \alpha_{k}^{\dagger} \alpha_{k}). \tag{9}$$

(b) Argue that a general eigenstate of the Heisenberg antiferromagnet in the spin-wave approximation can be written (we omit normalization factors)

$$\prod_{\mathbf{k}} [(\alpha_{\mathbf{k}}^{\dagger})^{n_{\alpha_{\mathbf{k}}}} (\beta_{\mathbf{k}}^{\dagger})^{n_{\beta_{\mathbf{k}}}}] |G\rangle \tag{10}$$

where  $\{n_{\alpha_k}\}$  and  $\{n_{\beta_k}\}$  are the magnon numbers characterizing the eigenstate and  $|G\rangle$  is the ground state discussed in the lectures. Show that this state is also an eigenstate of  $S^z$  and determine the corresponding eigenvalue.